

# Cosmology with Clusters

Alain Blanchard

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but  $\delta h \ll 1$  so Newton dynamics is enough.

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CMB:

$$\frac{\delta T}{T} \sim \delta h \sim \frac{\sigma^2}{c^2} \sim 10^{-5}$$

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but by CMB!

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Minimal assumption: gravity should be active.

# The Coma cluster

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# 3D surveys

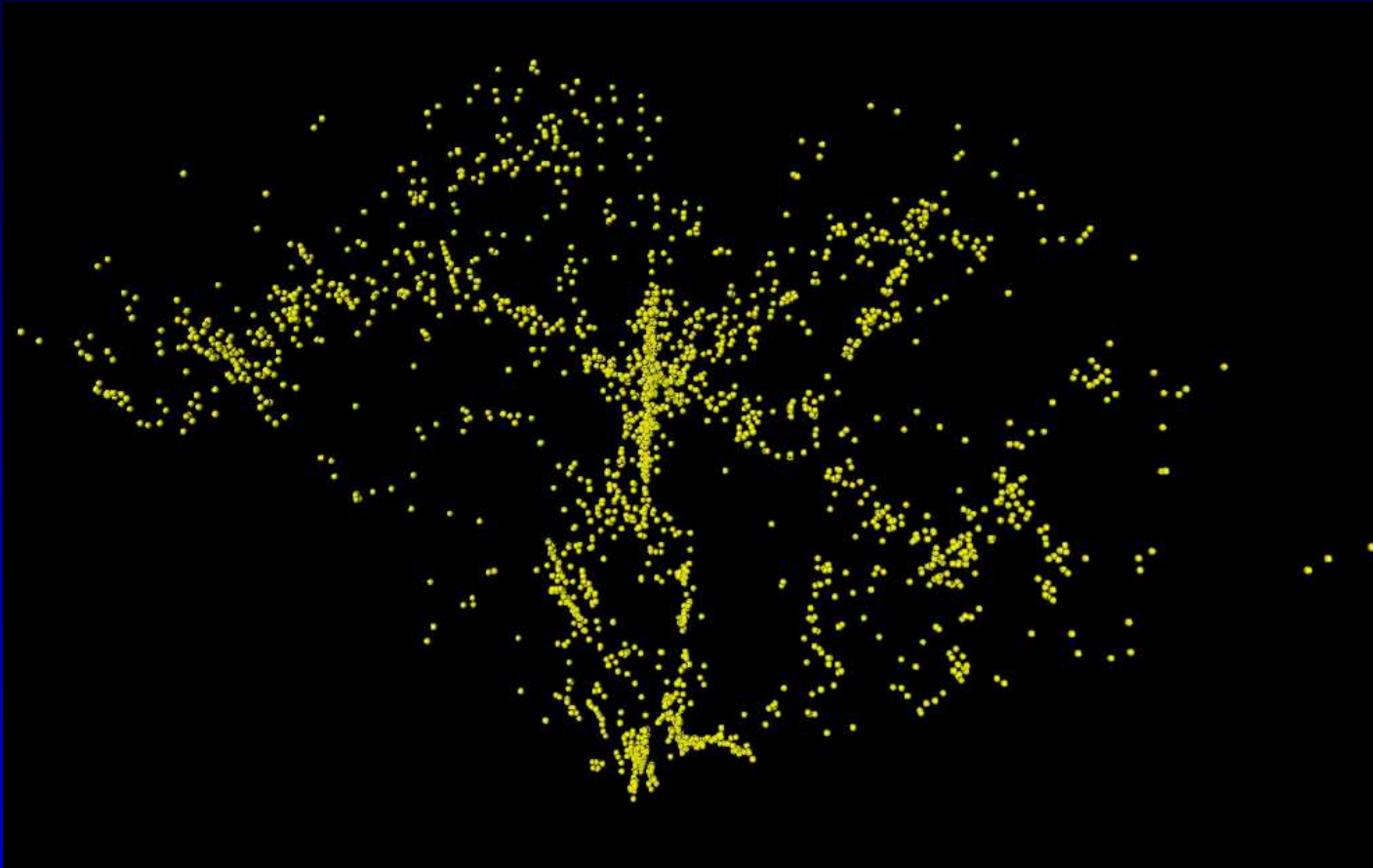
# 3D surveys

Velocity dispersion in galaxy clusters.



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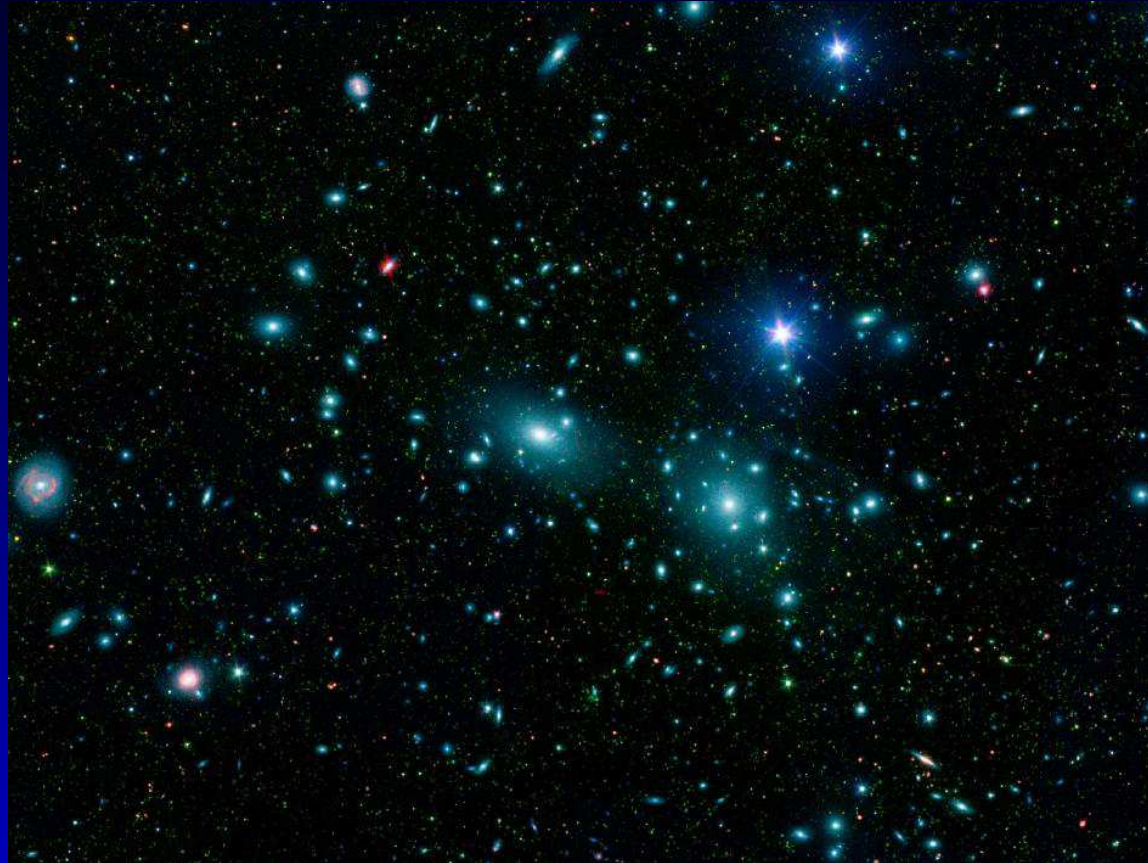
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Zwicky ( $\sim 1930$ ) inferred the presence of dark matter.

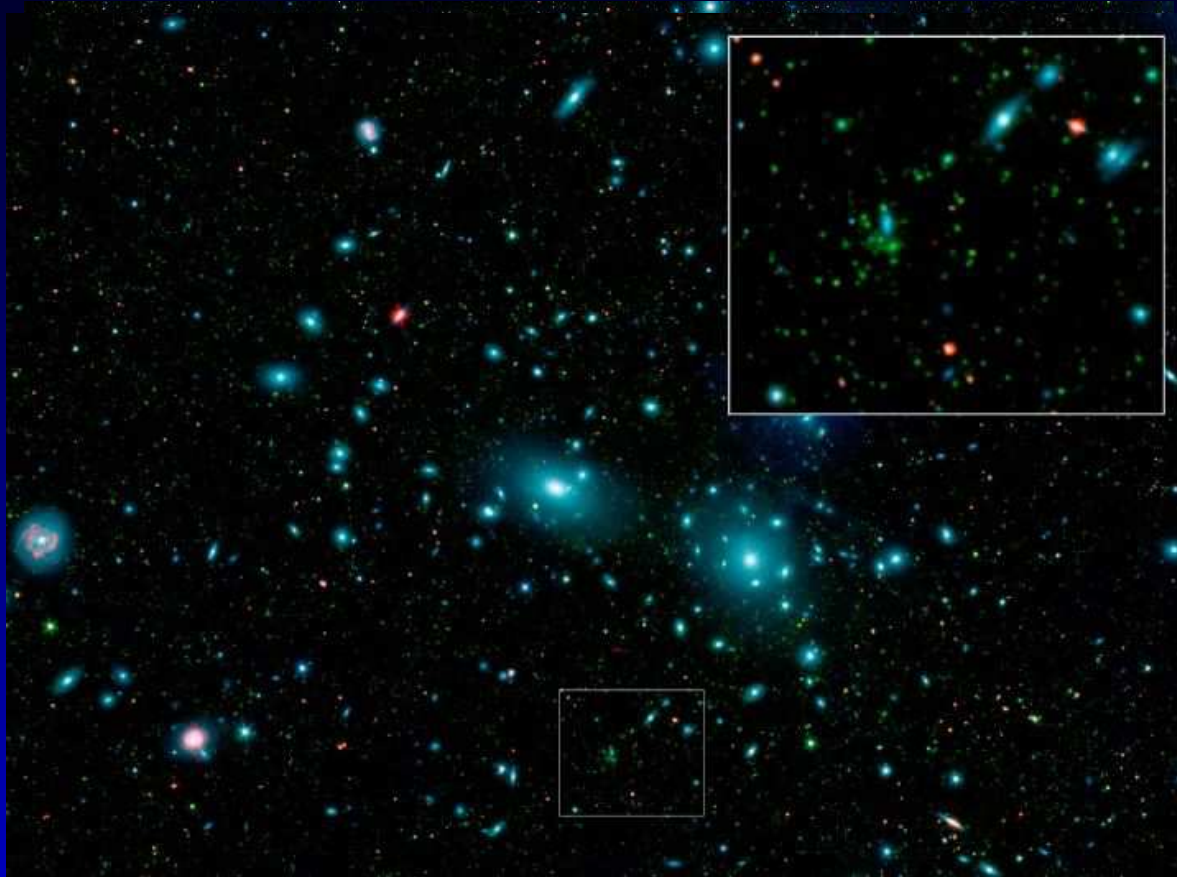
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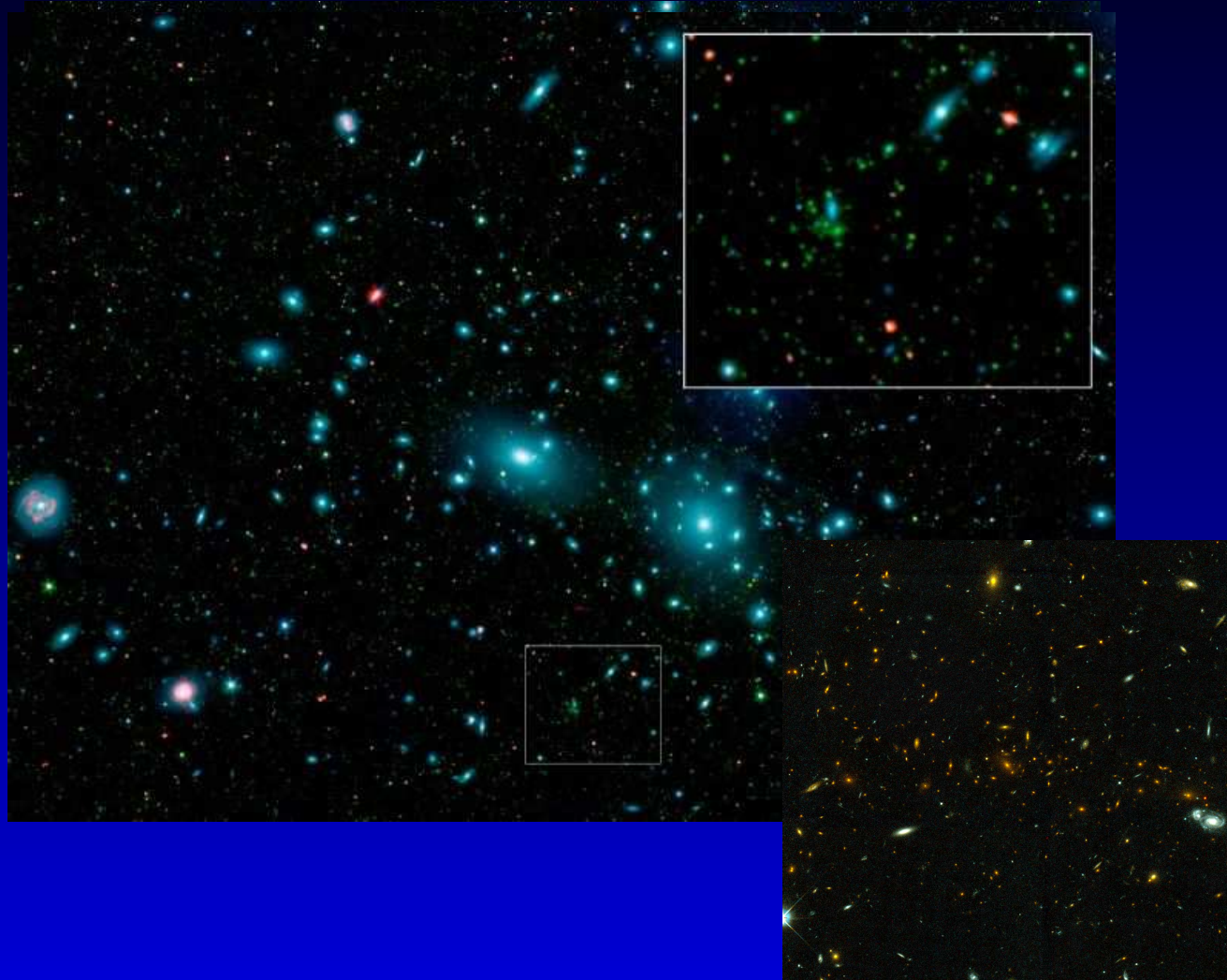




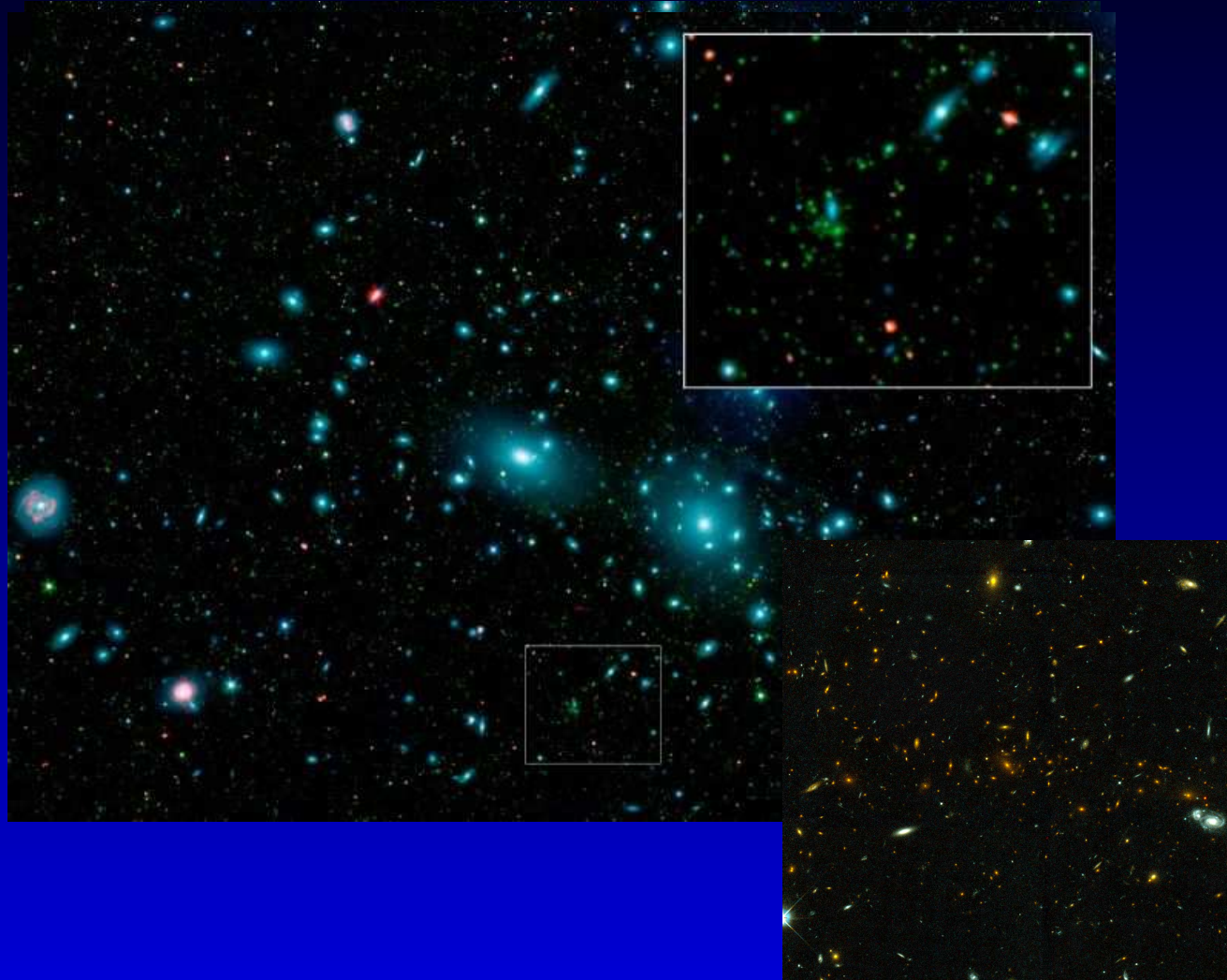
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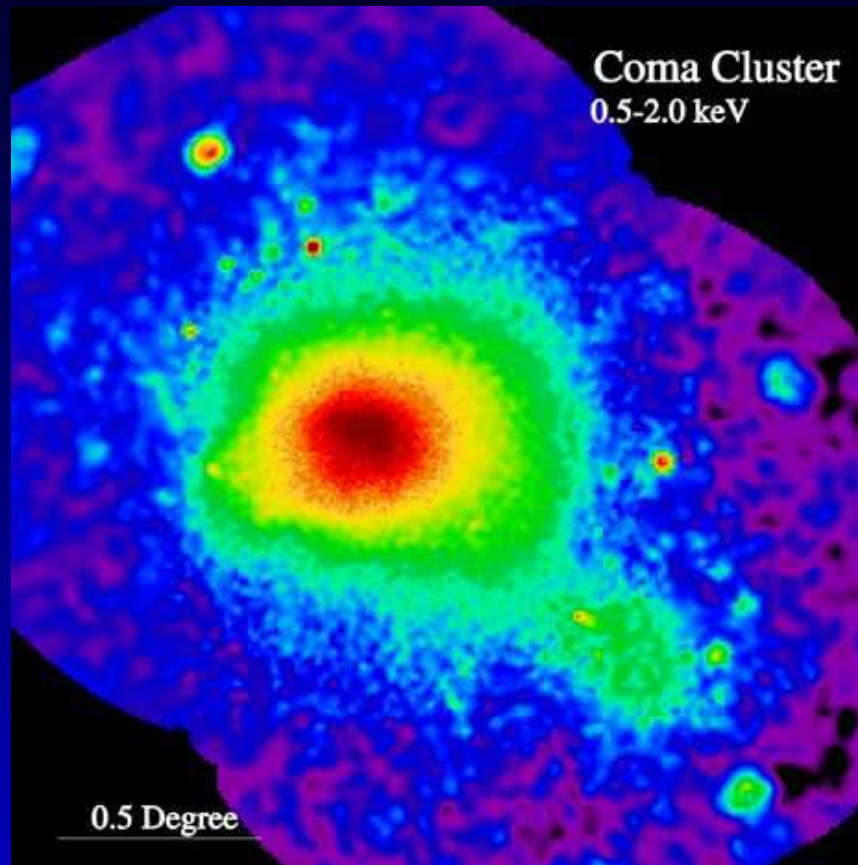


**Optical data** : Stars, metals, velocity dispersion →  
Mass...

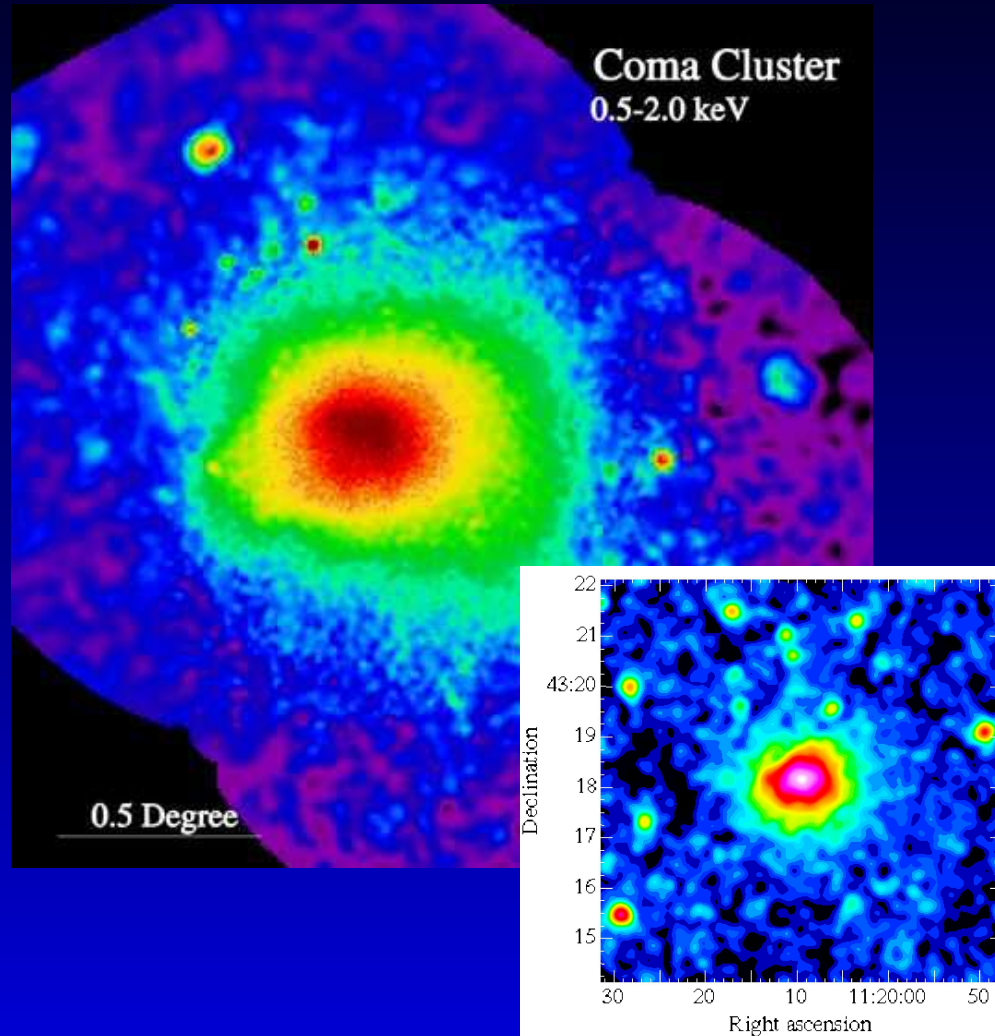
# X-ray Visions on clusters



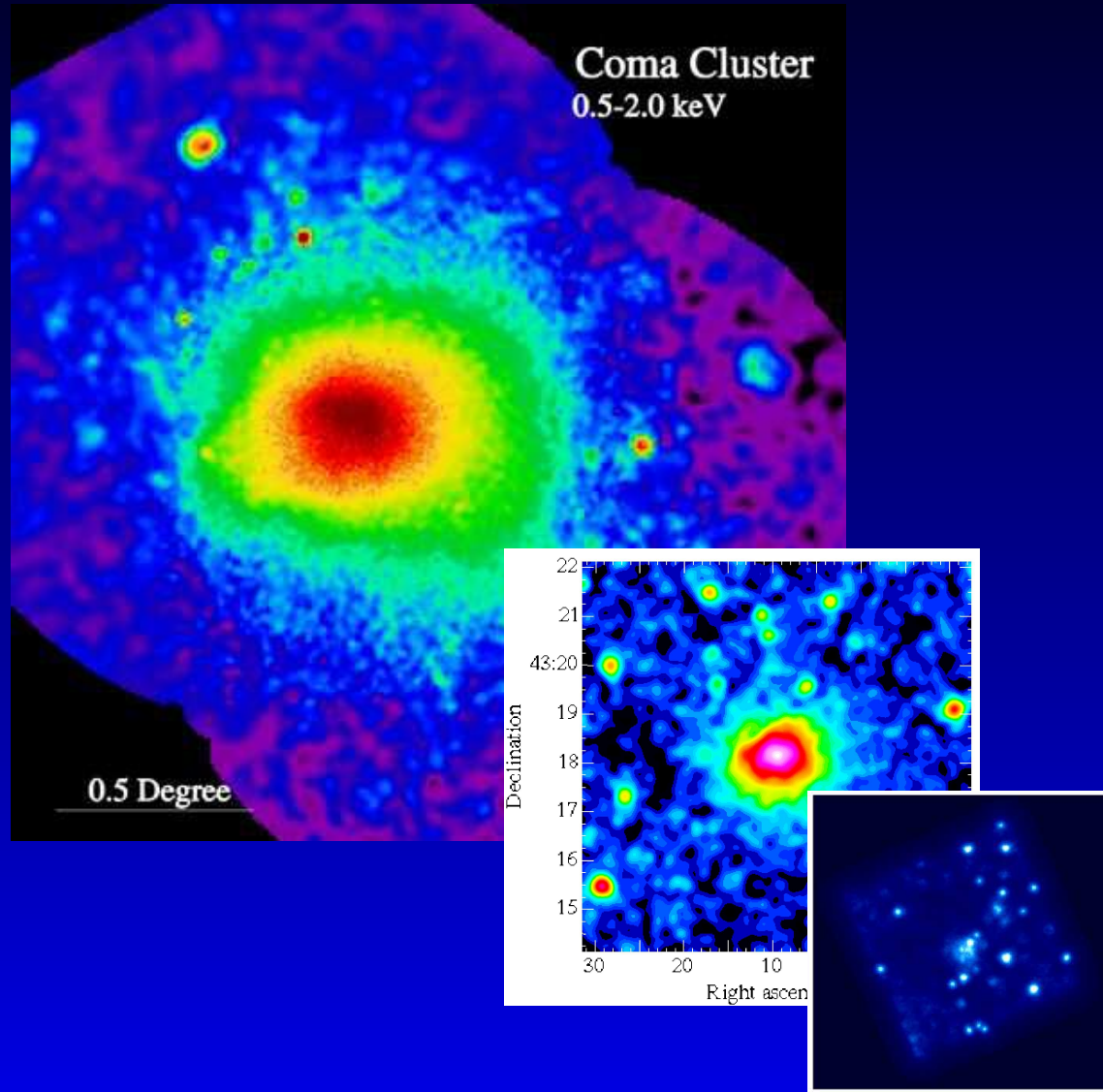
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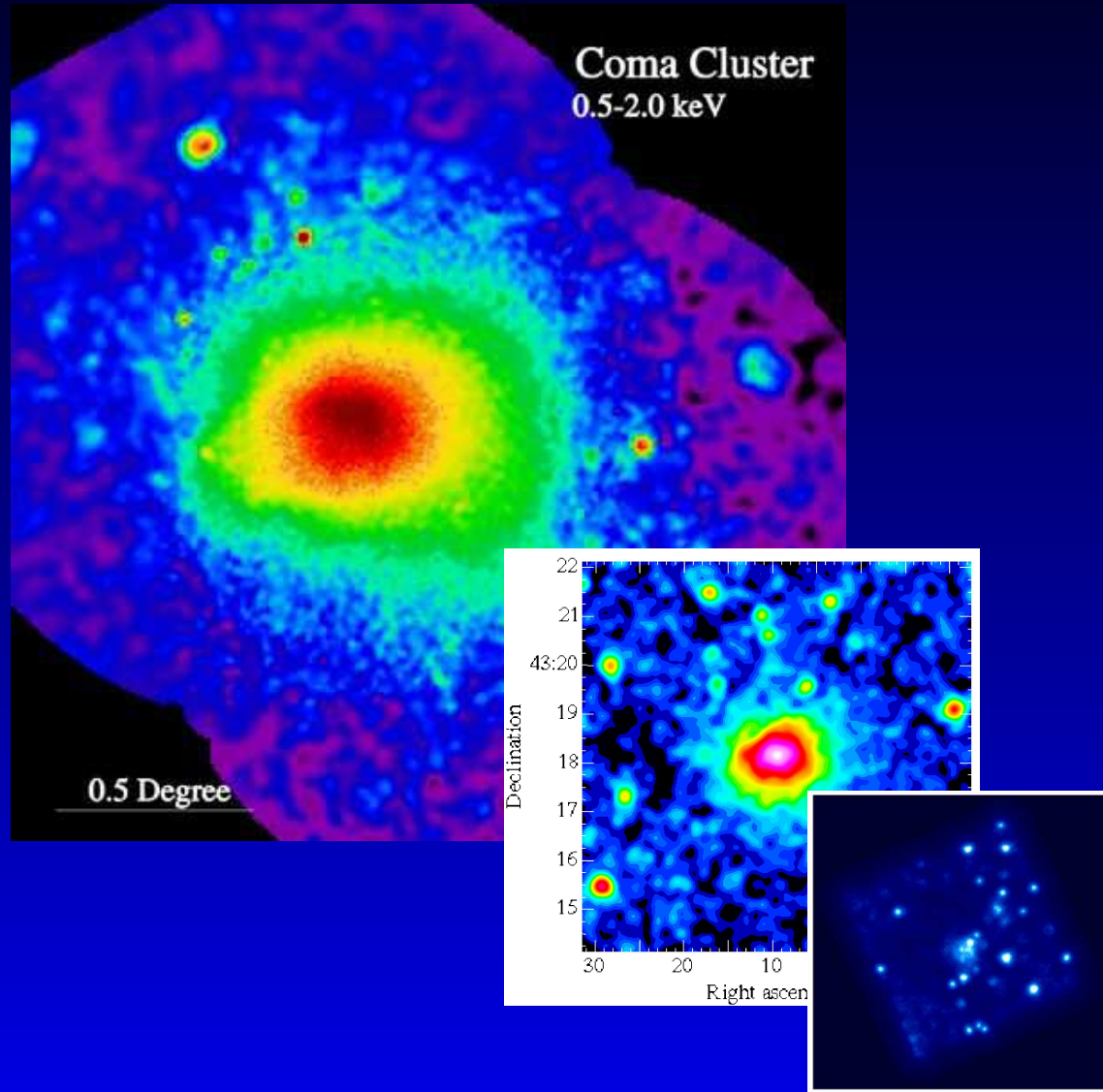
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**X-ray data : Gas, metals, temperature  $\rightarrow$  Mass...**

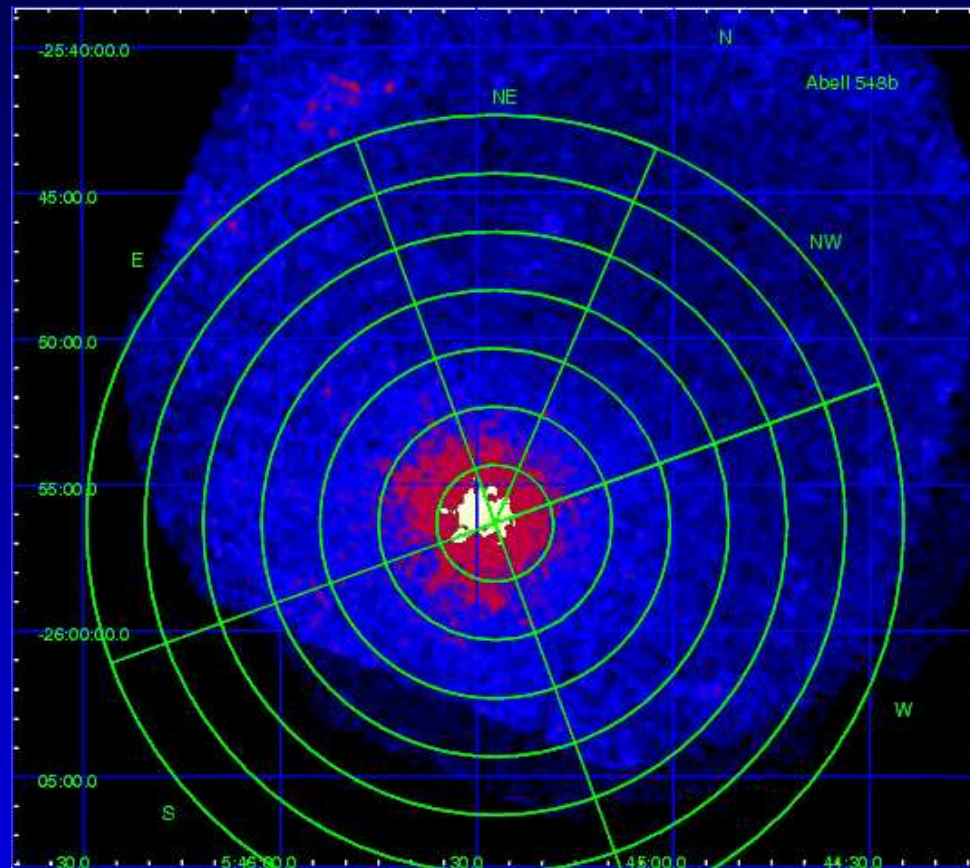


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XMM view of A548b

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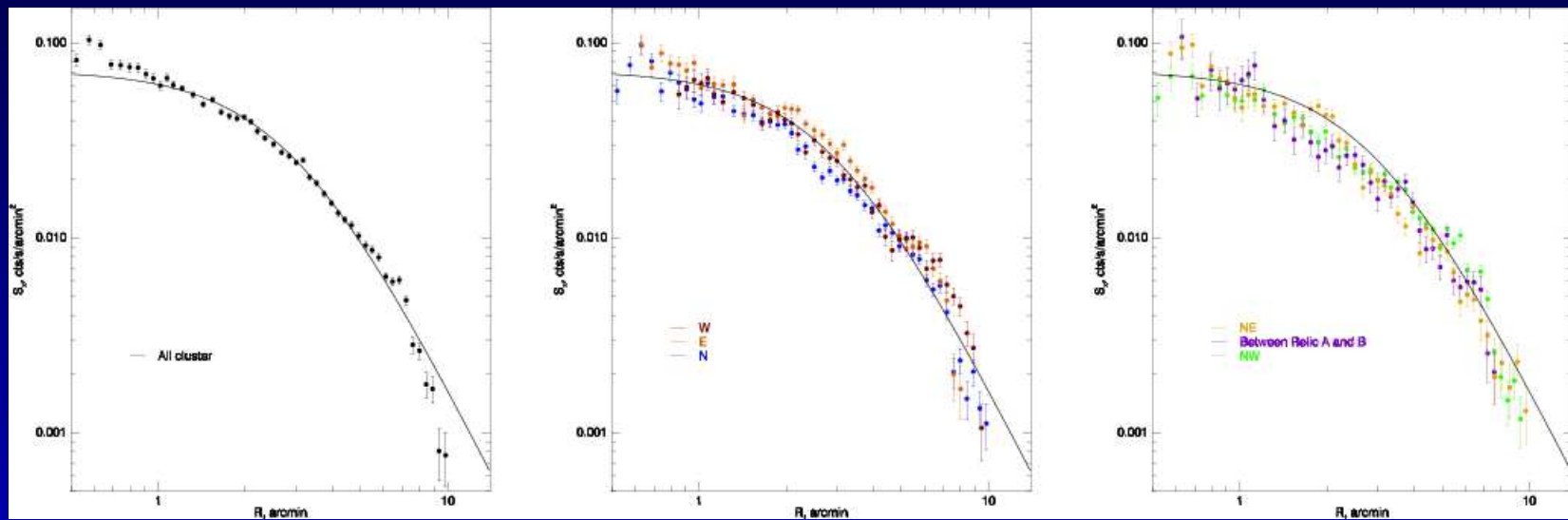


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XMM view of A548b: Sx profile

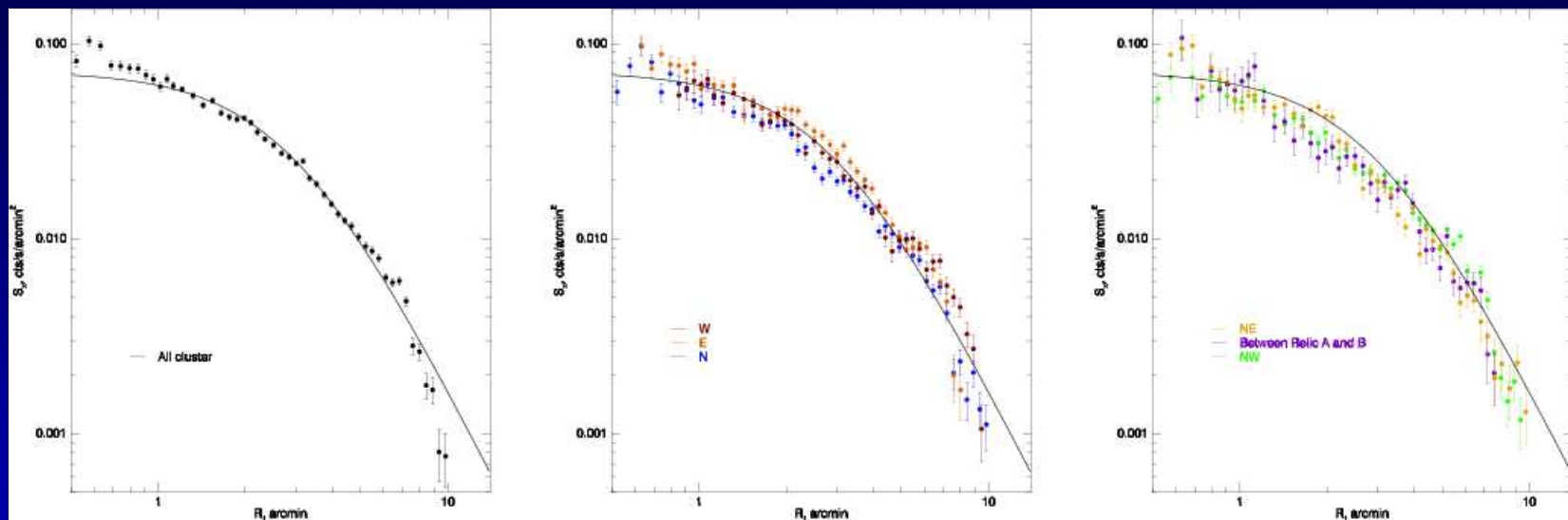
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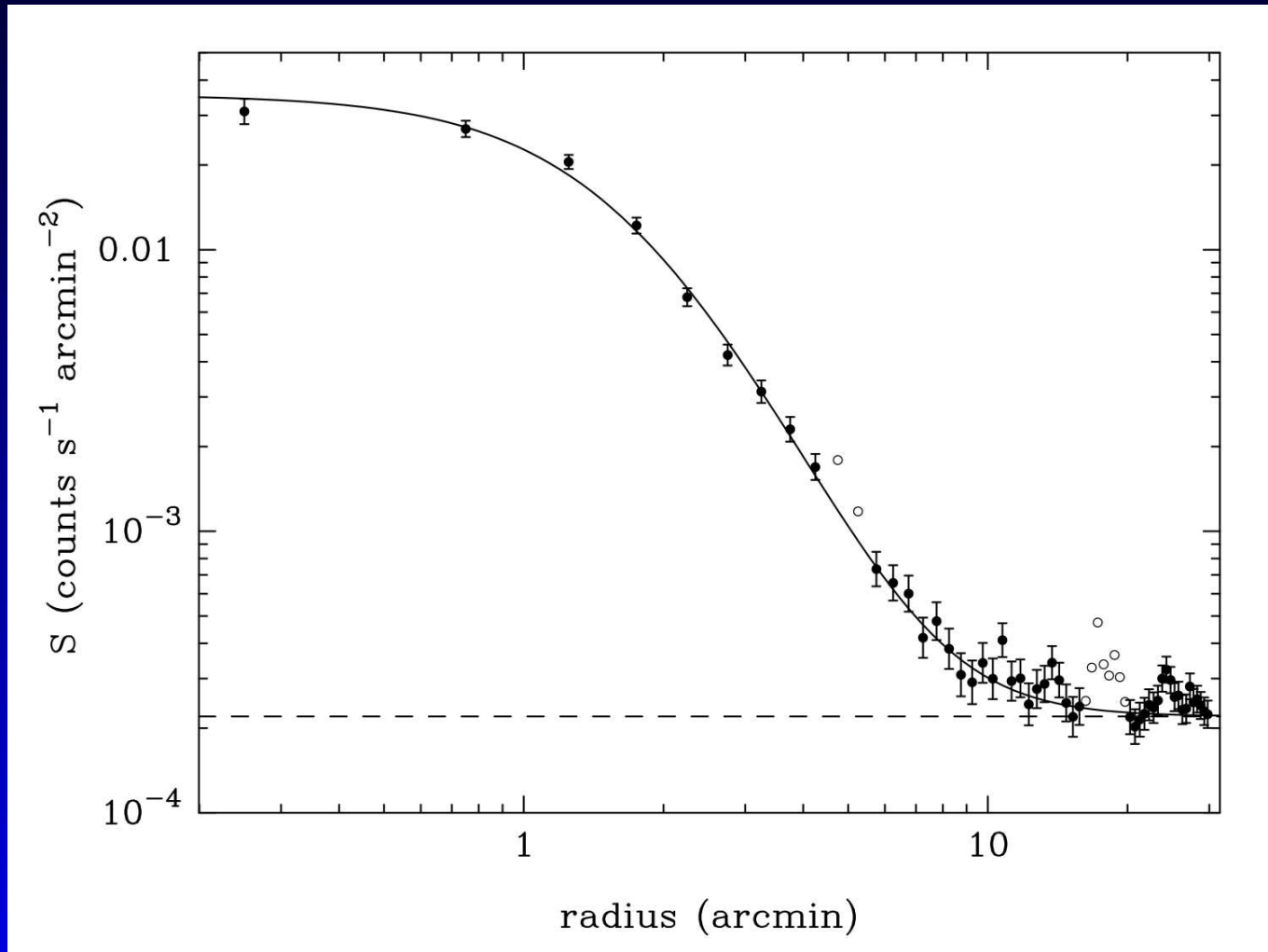
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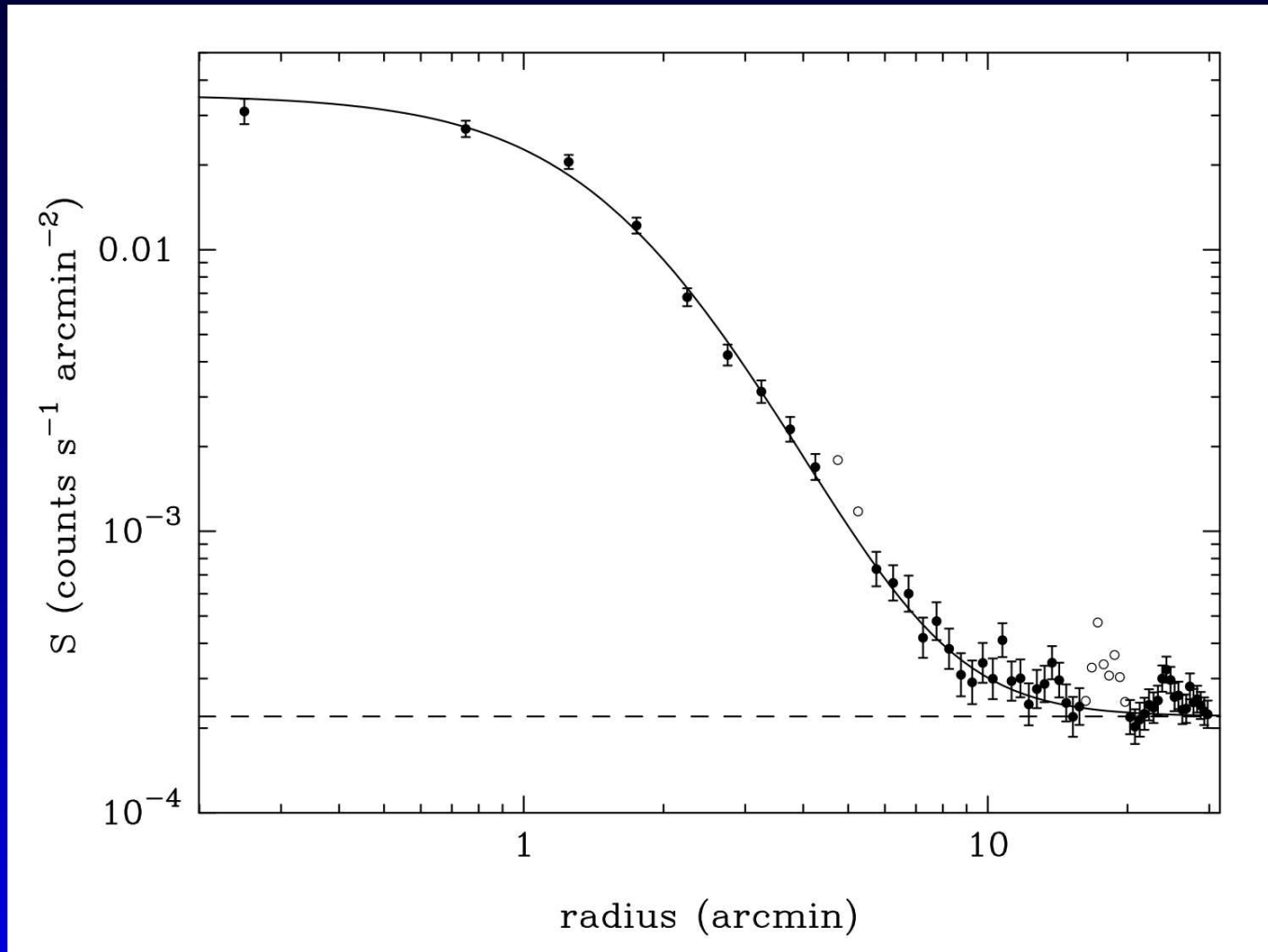
$$s(\theta) = \frac{s_0}{\left(1 + (\theta/\theta_c)^2\right)^{3\beta - 0.5}}$$

# Sx profile

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background has to be subtracted.

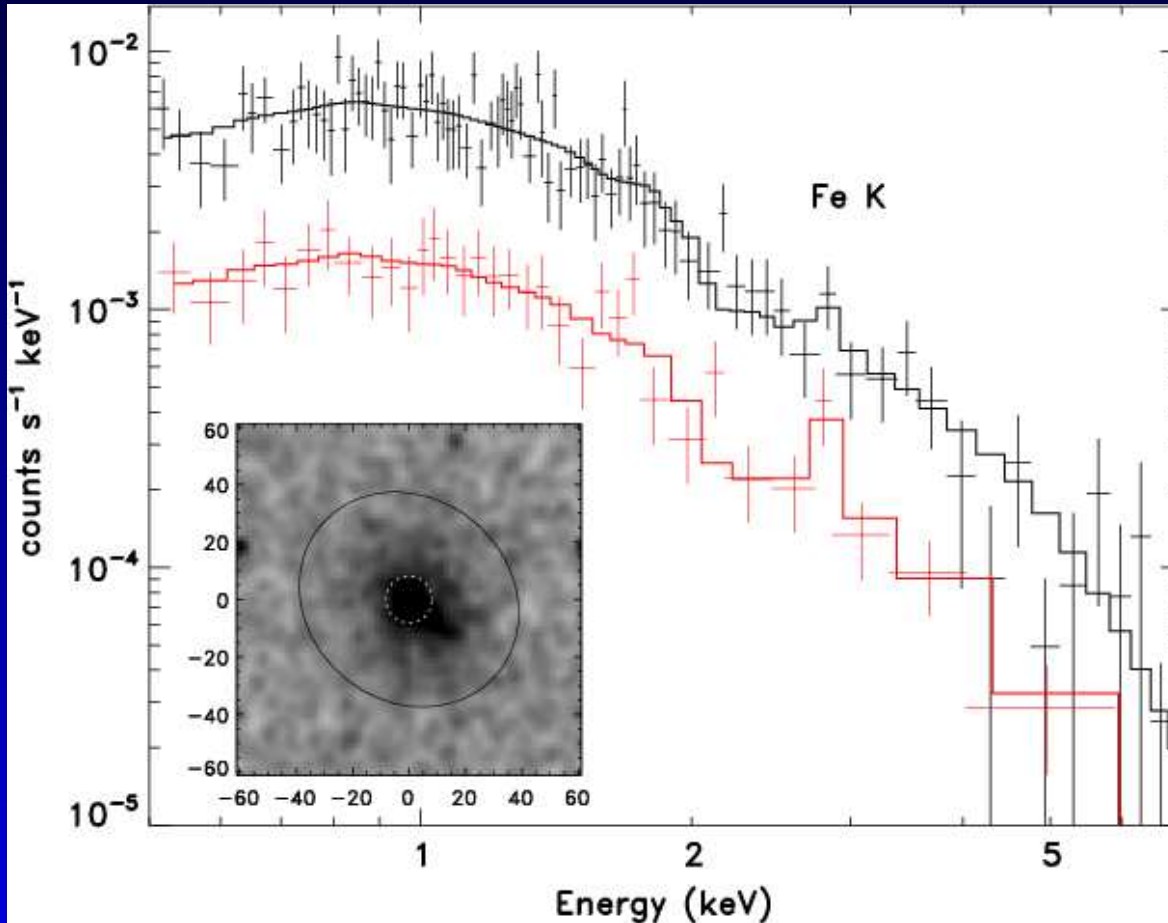


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# Hydrostatic mass estimates

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Isothermal beta model ( $r \gg r_c$ ) :

$$M(r) = -3\beta \frac{kT(r)}{G\mu m_p} r$$

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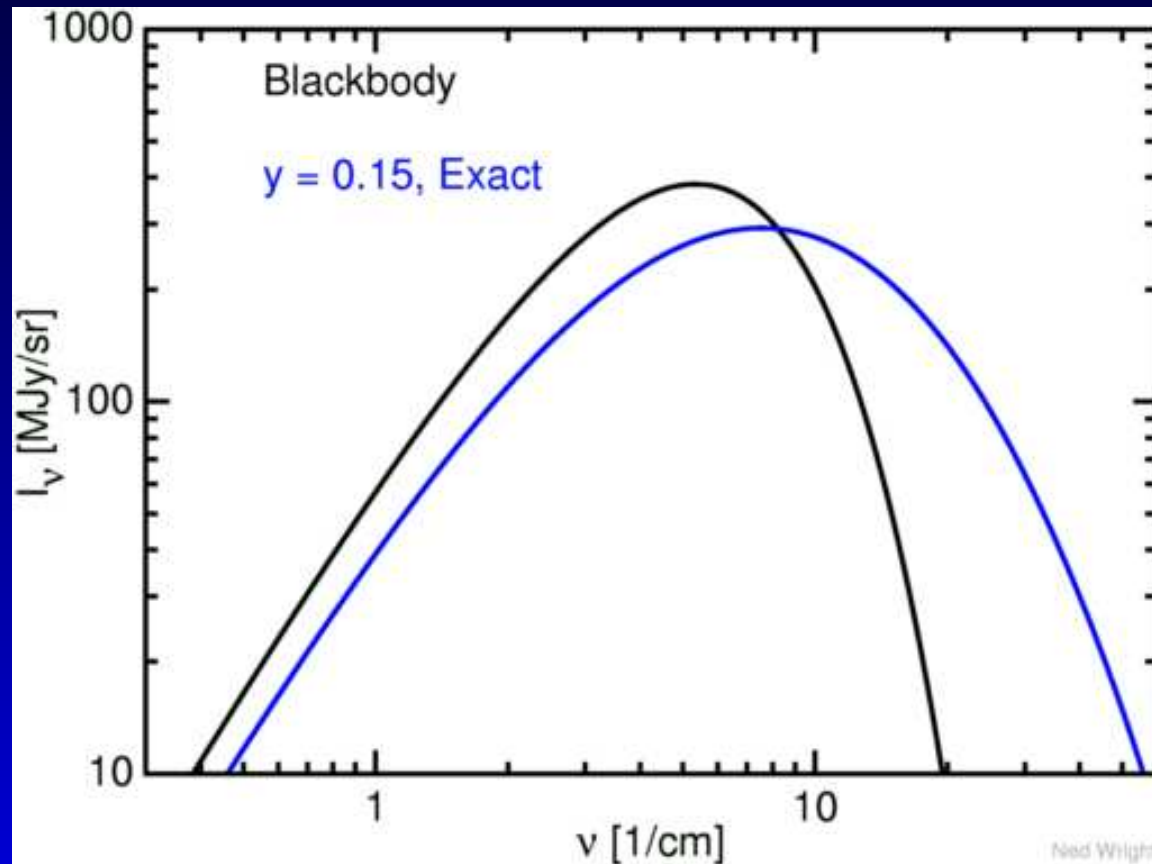
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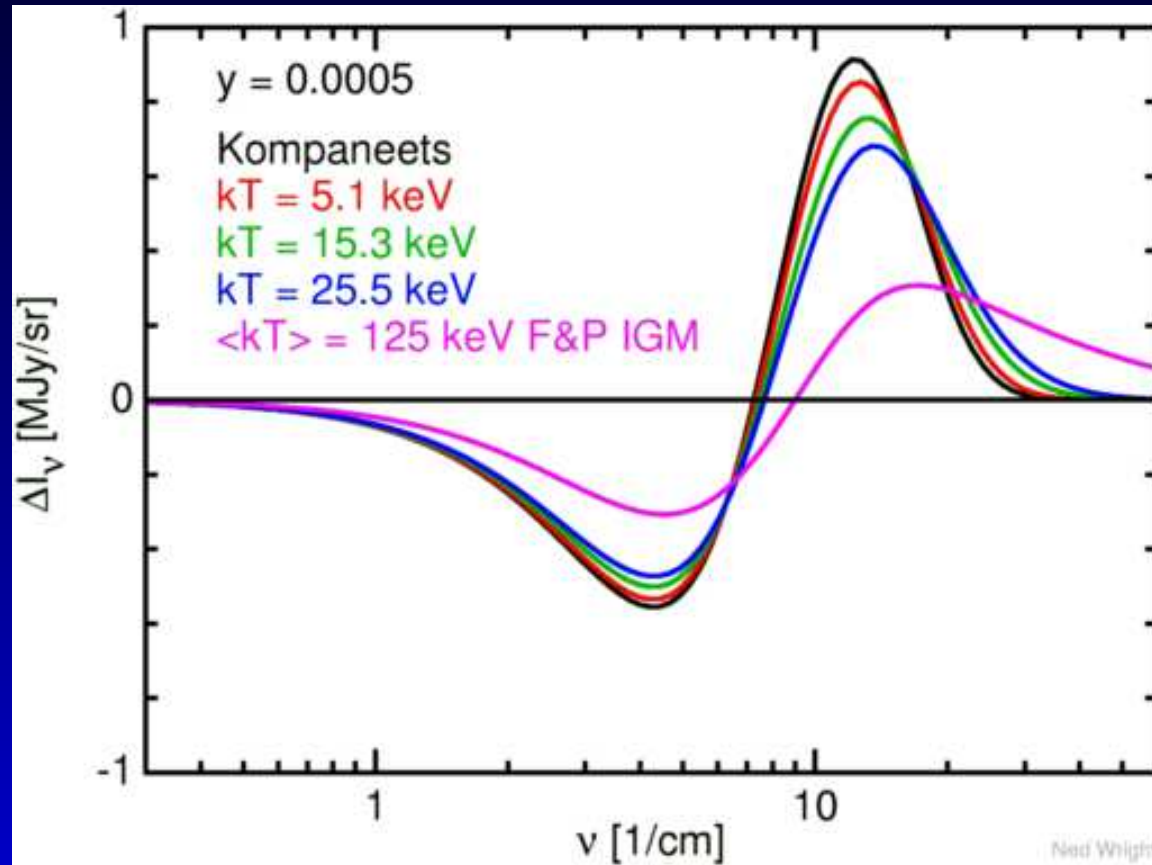
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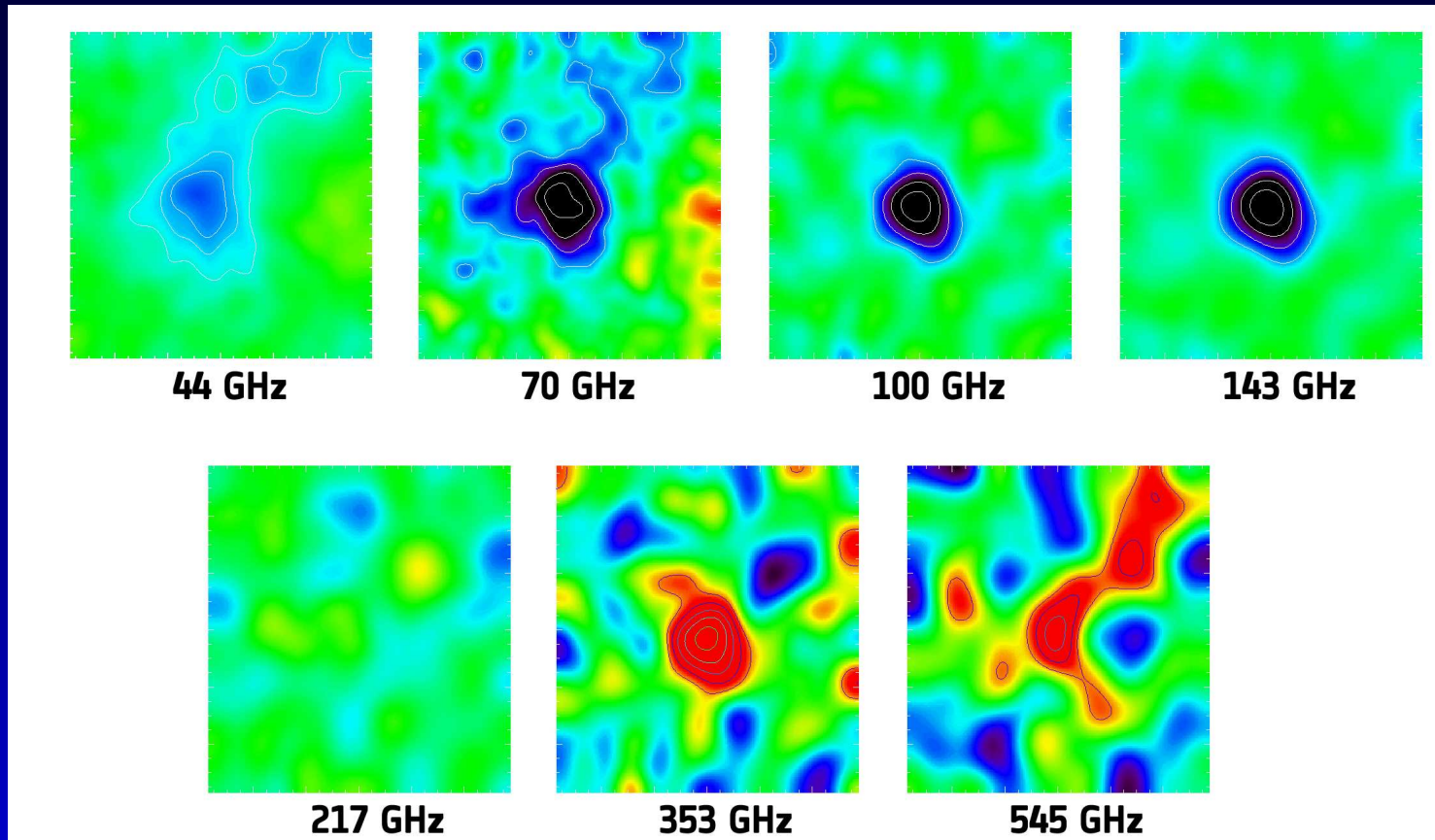
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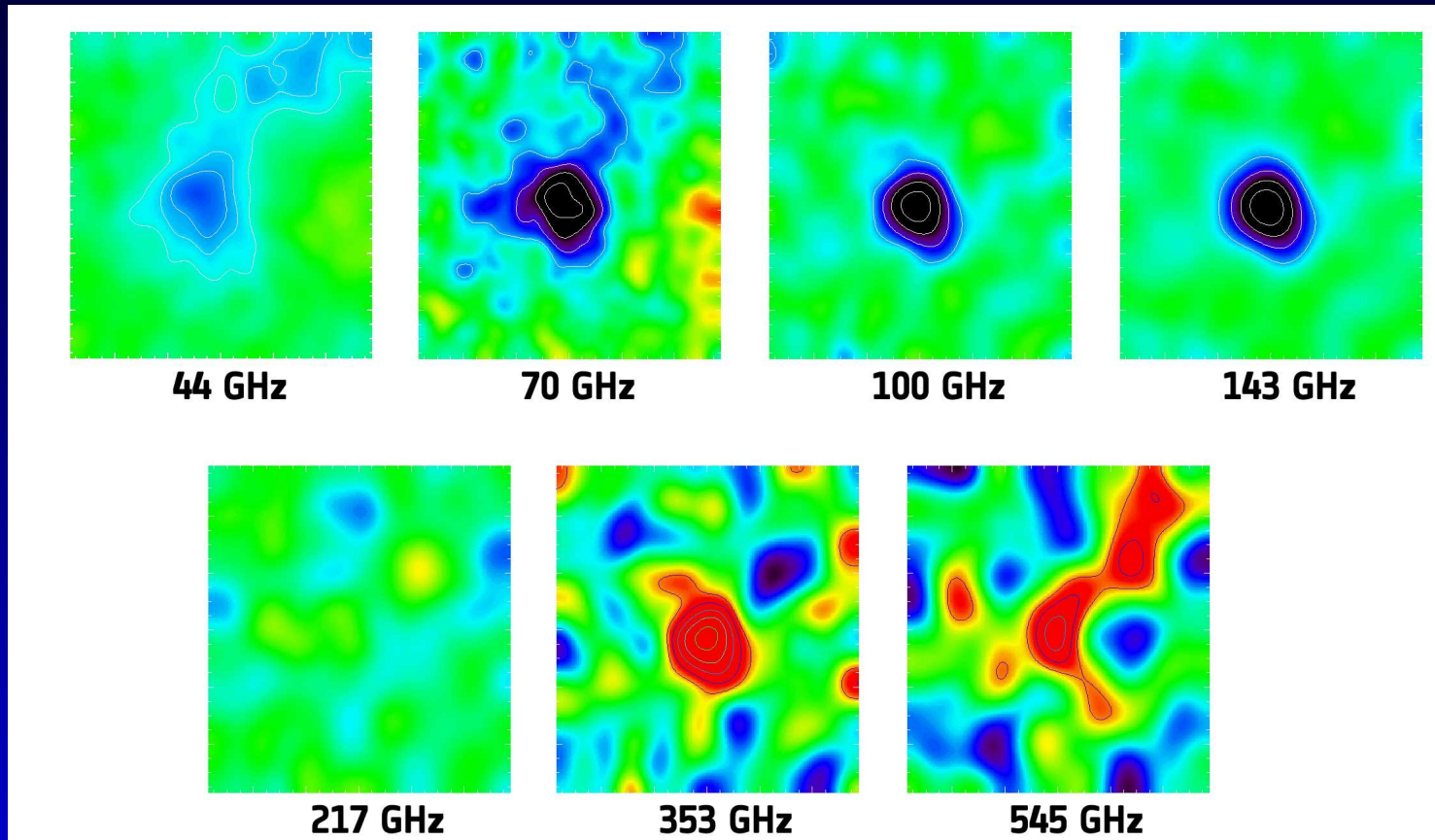
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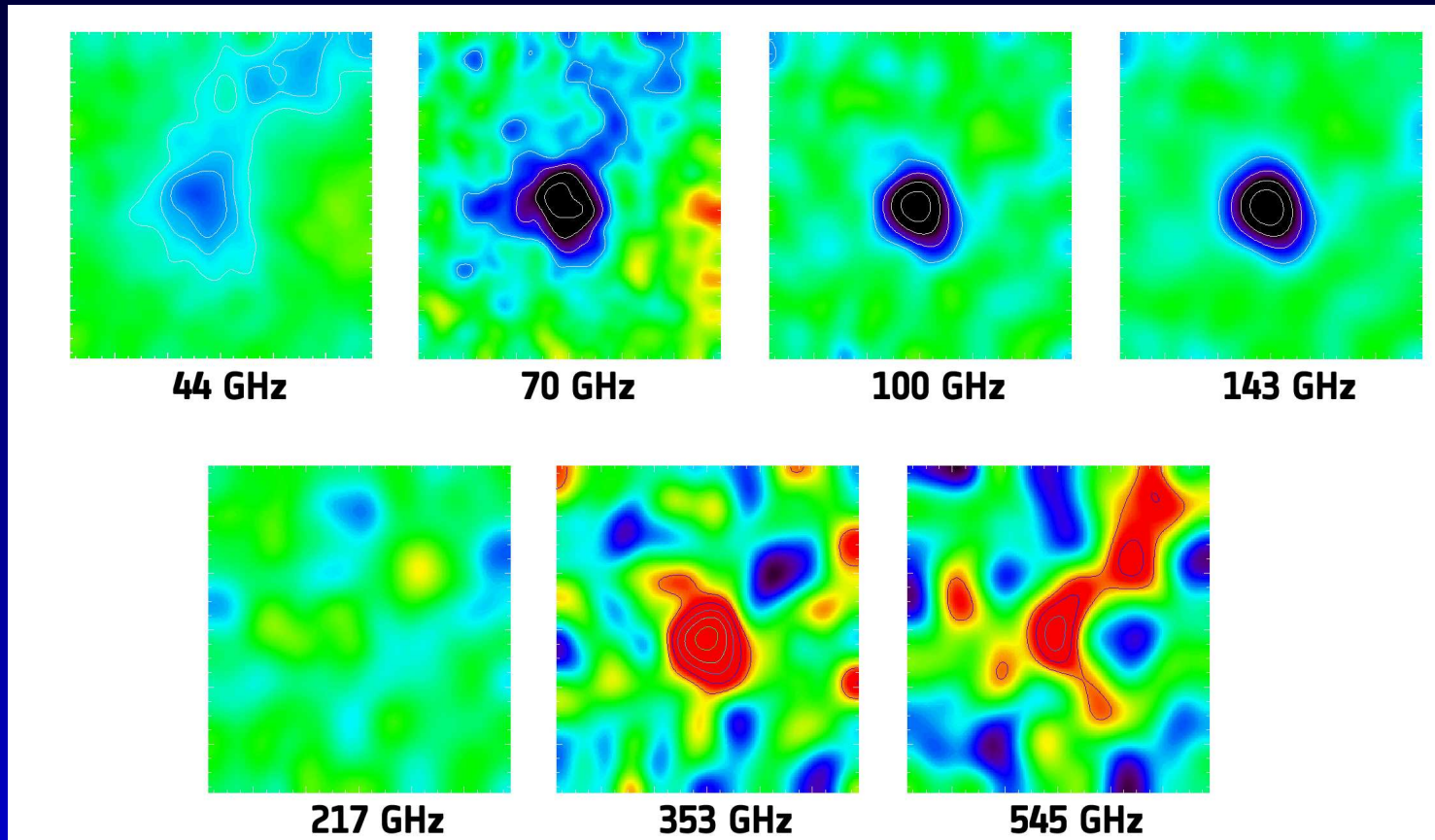


**SZ Signal** : Gas mass  $\times$  temperature  $\rightarrow$  Mass...



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No dimming with redshift



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- → fundamental probes for cosmology

# Non linear regime



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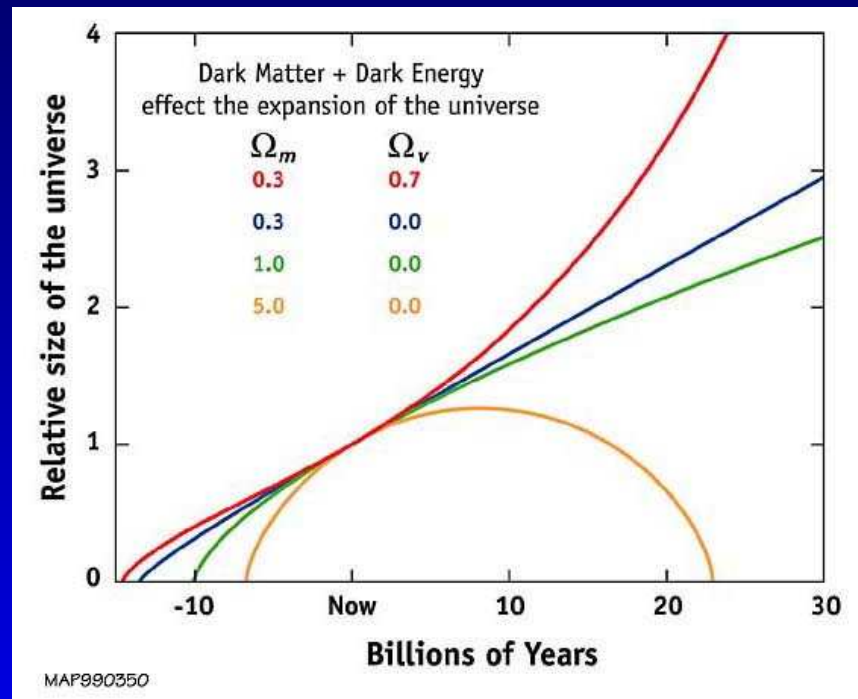
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Newtonian problem. Solution already seen:



# Spherical Perturbation I

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$$\tilde{H}_0 t = \frac{\tilde{\Omega}_0}{2(\tilde{\Omega}_0 - 1)^{3/2}} (\phi - \sin(\phi))$$
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Density at maximum:

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$$\Delta_m = \frac{9}{16} \pi^2 - 1. \simeq 4.55$$

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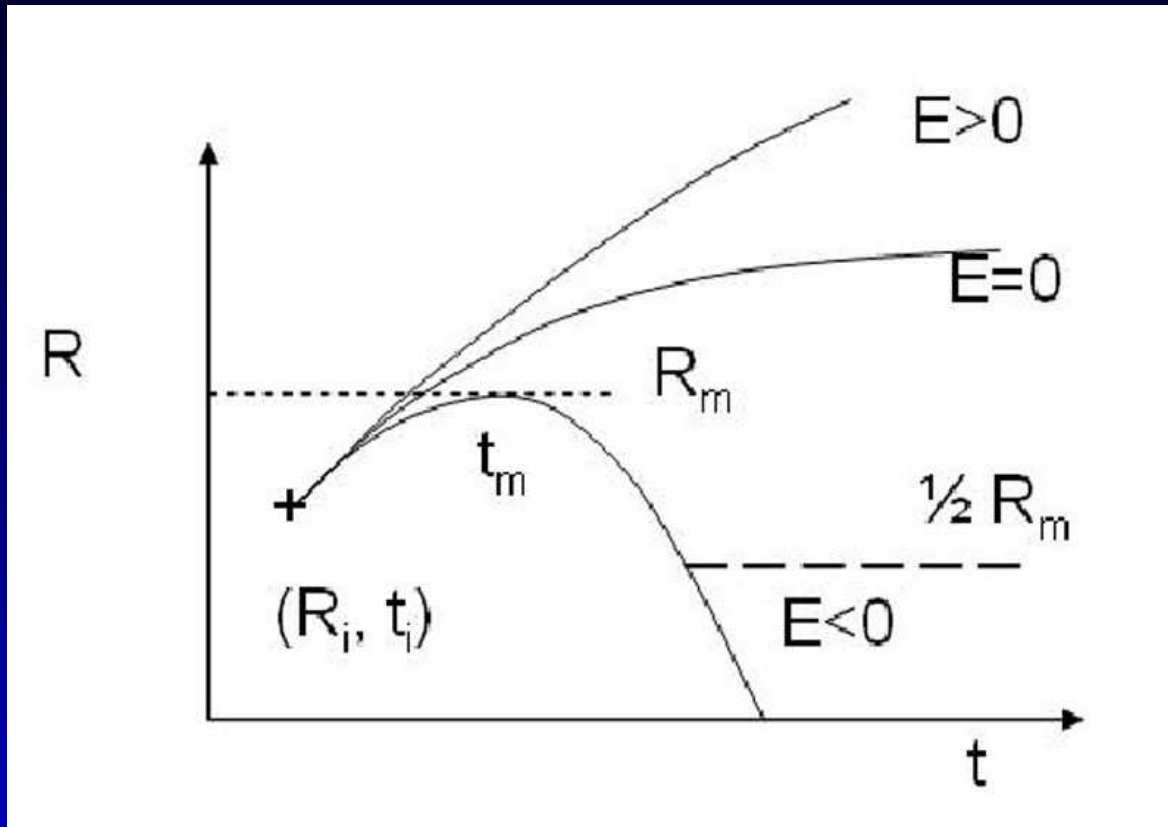
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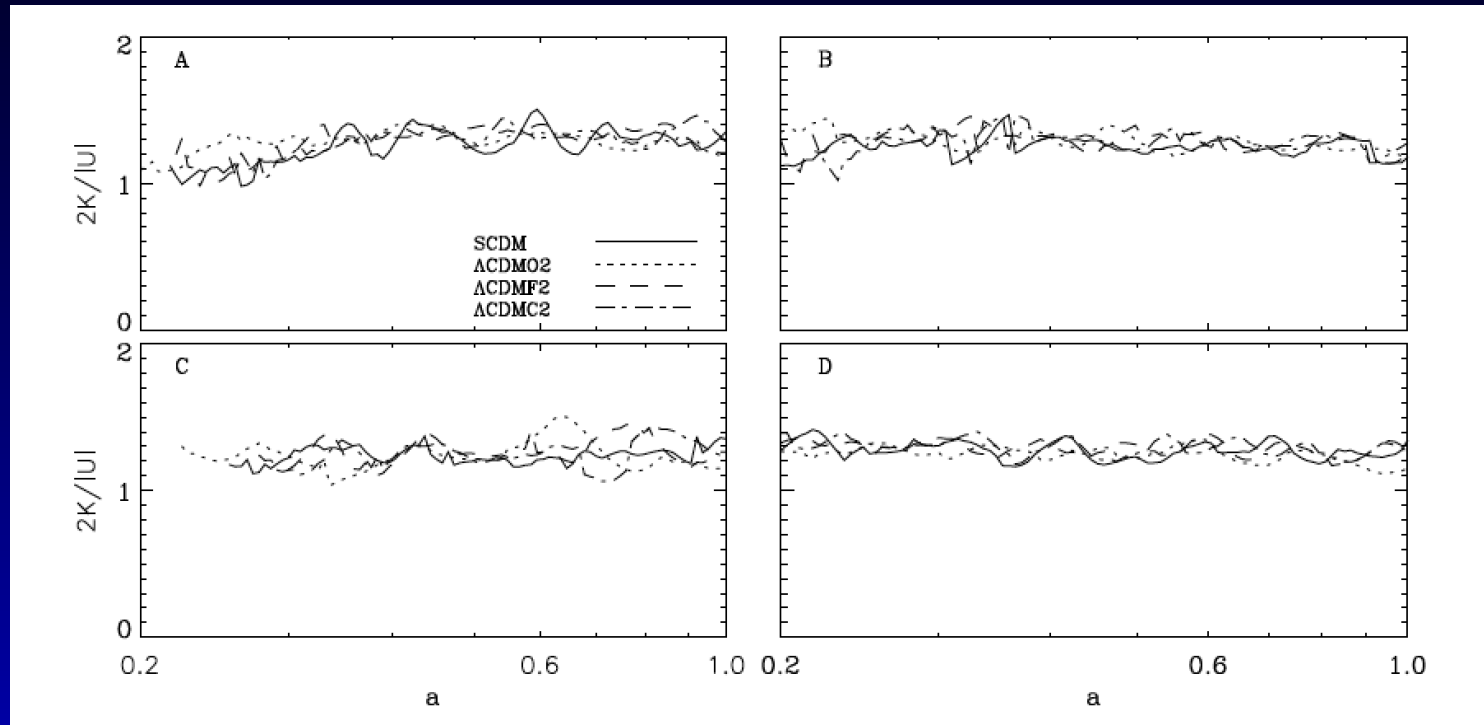
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let's estimate the linear expected amplitude at virilization.

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$$t = \frac{t_m}{\pi}(\psi - \sin \psi) = \frac{t_m}{\pi} \frac{\psi^3}{6} \left[1 - \frac{\psi^2}{20}\right]$$

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SO:

$$\psi^6 = \left(\frac{6\pi t}{t_m}\right)^2 \left[1 + \frac{\psi^2}{10}\right]$$



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and

$$\begin{aligned}\tilde{\rho} &= \left(1 + \frac{\psi^2}{4} - \frac{\psi^2}{10}\right) \frac{64\rho_m t_m^2}{(6\pi)^2 t^2} \\ &= \left(1 + \frac{\psi^2}{4} - \frac{\psi^2}{10}\right) \frac{64}{36\pi^2} \frac{3\pi^2}{32\pi G t^2} = \rho \left(1 + \frac{3\psi^2}{20}\right)\end{aligned}$$

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so with :  $\tilde{\rho} = \rho(1 + \delta)$

$$\delta = \frac{3}{20}\psi^2 = \frac{3}{20} \left(\frac{6\pi t}{t_m}\right)^{2/3} = \frac{3(6\pi)^{2/3}}{20} \frac{1 + z_m}{1 + z}$$

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$$\delta_m = \frac{3(6\pi)^{2/3}}{20}(1+z_m) = 1.06(1+z_m) \text{ when } \Delta_m \simeq 4.5$$

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$$\text{For } z < z_v, \Delta = 177 \left( \frac{1+z_v}{1+z} \right)^3$$



# Flat cosmology

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$$\Delta_v \approx 18\pi^2 + 82x - 39x^2$$

(compared to  $\rho_c(z)$ ) with  $x = \Omega(z) - 1$  (Bryan and Norman, 1998) and

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(for the spherical collapse model)

$$\Delta_v \approx 18\pi^2 + 82x - 39x^2$$

(compared to  $\rho_c(z)$ ) with  $x = \Omega(z) - 1$  (Bryan and Norman, 1998) and

$$\delta_{th} = \frac{3}{20} (12\pi)^{2/3} [1 + 0.0123 \log(\Omega(z))].$$

(Kitayama and Suto, 1996)